AIP The Journal of Chemical Physics

Origin of scaling behavior of protein packing density: A sequential Monte Carlo study of compact long chain polymers

Jinfeng Zhang, Rong Chen, Chao Tang, and Jie Liang

Citation: J. Chem. Phys. **118**, 6102 (2003); doi: 10.1063/1.1554395 View online: http://dx.doi.org/10.1063/1.1554395 View Table of Contents: http://jcp.aip.org/resource/1/JCPSA6/v118/i13 Published by the American Institute of Physics.

Additional information on J. Chem. Phys.

Journal Homepage: http://jcp.aip.org/ Journal Information: http://jcp.aip.org/about/about_the_journal Top downloads: http://jcp.aip.org/features/most_downloaded Information for Authors: http://jcp.aip.org/authors

ADVERTISEMENT



Origin of scaling behavior of protein packing density: A sequential Monte Carlo study of compact long chain polymers

Jinfeng Zhang

Department of Bioengineering, University of Illinois at Chicago, Chicago, Illinois 60607

Rong Chen

Departments of Information and Decision Science and Bioengineering, University of Illinois at Chicago, Chicago, Illinois 60607

Chao Tang

NEC Research Institute, Princeton, New Jersey 08540

Jie Liang^{a)}

Department of Bioengineering, University of Illinois at Chicago, Chicago, Illinois 60607

(Received 14 October 2002; accepted 31 December 2002)

Single domain proteins are thought to be tightly packed. The introduction of voids by mutations is often regarded as destabilizing. In this study we show that packing density for single domain proteins decreases with chain length. We find that the radius of gyration provides a poor description of protein packing but the alpha contact number we introduce here characterize proteins well. We further demonstrate that protein-like scaling relationship between packing density and chain length is observed in off-lattice self-avoiding walks. A key problem in studying compact chain polymers is the attrition problem: It is difficult to generate independent samples of compact long self-avoiding walks. We develop an algorithm based on the framework of sequential Monte Carlo and succeed in generating populations of compact long chain off-lattice polymers up to length N=2000. Results based on analysis of these chain polymers suggest that maintaining high packing density with chain length of proteins is a generic feature of random polymers satisfying loose constraint in compactness. We conclude that proteins are not optimized by evolution to eliminate packing voids. © 2003 American Institute of Physics. [DOI: 10.1063/1.1554395]

I. INTRODUCTION

Geometric considerations have lead to important insights about protein structures.¹⁻⁵ Voids are simple geometric features that represent packing defects inside protein structures. For multisubunit proteins such as GroEL and potassium channel, voids or tunnels of large size are formed by the spatial arrangement of multiple subunits, and are essential for the biological functions of these proteins.^{6,7} In this study, we focus on voids formed due to packing defects that are not directly involved in protein function. For this purpose, we choose to study only structures of single domain proteins. Although these proteins are well known to be compact,⁸ and their interior is frequently thought to be solid-like,^{9,10} recent calculations showed that there are also numerous voids buried in the protein interior.¹¹ The importance of tight packing in single chain protein is widely appreciated; packing is thought to be important for protein stability,¹²⁻¹⁴ for kinetic nucleation of protein folding,^{15,16} and for successful design of novel proteins following a predefined backbone.¹⁴ The conservation of amino acid residues during evolution may also be correlated with tightly packed sites.^{15–17} In contrast, the potential roles of voids in affecting protein stability and in influencing tolerance to mutations and designability of proteins^{18,19} are not well understood.

An important parameter describing packing is the packing density p_d , which is a quantitative measure of the voids and was first introduced to study proteins by structural biologists. This concept has been widely used in protein chemistry.^{8,13} The scaling relationship of p_d and chain length N was first studied in Ref. 11. p_d can be thought of as the physical volume v_{vdw} occupied by the union of van der Waals atoms, divided by the volume of an envelope v_{env} that tightly wraps around the body of atoms, $p_d \equiv v_{vdw}/v_{env}$.¹¹ Voids contained within the molecule will not be part of the van der Waals volume v_{vdw} , but will be included in v_{env} . Using geometric algorithms, v_{vdw} , v_{env} , and p_d can be readily computed for protein structures in the Protein Data Bank.^{20,21}

In this work, we further study the scaling behavior of packing density p_d with chain length of single domain proteins and explore the determinants of the observed scaling behavior. We seek to answer the following questions: Is the scaling behavior of p_d unique to proteins? Are proteins optimized during evolution to eliminate packing voids? We introduce two new packing parameters n_{α} (the alpha contact number) and z_{α} (the alpha coordination number). We show that n_{α} characterizes protein packing very well with a linear

^{a)}Author to whom correspondence should be addressed. Phone: (312)355-1789; Fax: (312)996-5921; Electronic mail: jliang@uic.edu

scaling relationship with the chain length, and that a widely used parameter, the radius of gyration R_g , characterizes protein packing poorly. To overcome the attrition problem of low success rate in generating compact long chain polymers, we develop an algorithm based on sequential Monte Carlo importance sampling and succeed in obtaining thousands of very compact long chain off-lattice polymers up to N= 2000. We demonstrate that the scaling behavior of p_d for proteins can be qualitatively reproduced by randomly generated polymers with rudimentary constraints of n_{α} . Our simulation studies lead us to conclude that proteins are not optimized to eliminate voids during evolution. Rather, voids in proteins are a generic feature of random polymers with a "reasonable" (as measure by z_{α}) compactness.

The paper is organized as follows: In the Methods, we first describe briefly how p_d and n_α are computed from the dual simplicial complex of protein structure, and introduce an off-lattice discrete model for generating random polymer conformations. We next describe the sequential Monte Carlo importance sampling and resampling techniques that allow us to generate adequate samples satisfying various criteria of n_α . In the Results, we begin with the characterization of void properties of proteins by both p_d and R_g . We then show the linear scaling behavior of n_α found in proteins. The scaling behavior of p_d of random polymers generated by sequential Monte Carlo with chain length is discussed later. We conclude with summary and discussion of our results.

II. METHODS

A. Protein data

To avoid complications of multichain and multidomain proteins, we examine the packing density of proteins of single domain proteins. We collect proteins from the PDBSELECT database²² that contains only one domain, as defined as single chains in the SCOP database with one numerical label.²³

B. Dual simplicial complex, alpha coordination number, and packing density

We use alpha shape to characterize the geometry of protein structure. Alpha shape has been successfully applied to study a number of problems in proteins, including void measurement, binding site characterization, protein packing, electrostatic calculations, and protein hydrations.^{11,20,21,24–28} Briefly, we first obtain a Delaunay simplicial complex of the molecule from weighted Delaunay triangulation, which decomposes the convex hull of atom centers into tetrahedra (3-simplices), triangles (2-simplices), edges (1-simplices), and vertices (0-simplices). We then obtain the dual simplicial complex of the protein molecule by removing any tetrahedra, triangles, and edges whose corresponding Voronoi vertices, edges, and planar facets are not at least partially contained within the protein molecule.^{29,30} The edges between atoms that are not connected by bonds corresponds to nonbonded alpha contacts. The total sum of the number of such edges for each atom is the total number of *alpha contacts* n_{α} . It reflects the total number of atoms that are in physical nearest neighbor contact with other atoms. These atoms have volume



FIG. 1. The alpha contacts in a toy molecule. In this molecule, both atom 1 and atom 2 have 4 alpha contacts. The number of atoms n=9, the number of alpha contacts is $n_{\alpha}=22$ (twice the number of edges), and the alpha coordination number $z_{\alpha}=n_{\alpha}/n\approx2.4$.

overlap and their corresponding weighted Voronoi cells intersect. The *alpha coordination number* is $z_{\alpha} \equiv n_{\alpha}/n$, where *n* is the total number of atoms in the molecule (see Fig. 1). In our calculation, we only consider nonbonded alpha contacts. Details of the theory and computation of alpha shape and dual simplicial complex can be found elsewhere.^{20,21,31}

We follow previous work in Ref. 11 and define packing density p_d as

$$p_d \equiv \frac{v_{\rm vdw}}{v_{\rm env}} = \frac{v_{\rm vdw}}{v_{\rm ms} + v_{\rm voids}},$$

where v_{vdw} , v_{ms} , and v_{voids} are van der Waals volume, molecular surface volume, and the void volume of the molecule, respectively.²⁰ Packing density is computed with a solvent probe radius 1.4 Å, as described in Ref. 11.

C. Growth model for off-lattice random polymers

We use a modified off-lattice discrete *m*-state model first developed in Ref. 32 to generate self-avoiding walks (SAWs) in three-dimensional space. All monomers are treated as balls with a radius of 1.7 Å. For monomers *i* and *j* that are not sequence near neighbors (|i-j|>2), the Euclidean distance d(i,j) between them must be greater than 2×1.7 Å so they are self-avoiding. Sequence neighboring monomers are connected by a bond of length 1.5 Å (Fig. 2).

We use a chain growth model to obtain conformation of polymer of specified length.³³ There are m=32 possible states where the next monomer can be placed. They are evenly distributed spatially on a sphere of radius 1.5 Å centered at the current monomer. We forbid the placement of the new monomer anywhere on a cap of the sphere with an angle $<60^{\circ}$ from the entering bond. This ensures that there are no unnatural acute sharp bond angles. The remaining sphere is divided into four strips, each may have different width but is of equal surface area. For the 32 possible states, we place uniformly 8 points at the midline of each of the 4 strips. Following Park and Levitt,³² the coordinates of each state are parametrized by two angles α and τ for ease of computation. α is the bond angle formed by the i-1, i and i+1 th monomers. τ is the torsion angle formed by four consecutive monomers.



FIG. 2. The 32-state discrete model for chain growth. There are 32 possible positions for adding the next monomer. They are located on a sphere of radius 1.5 Å, but placement on the cap with an angle $<60^{\circ}$ from the entering bond is forbidden. The surface is divided into four stripes of equal area. Eight positions are placed evenly on each strip. α is the bond angle.

D. Approximately maximum compact polymer

In addition, we generate polymers that are approximately maximum compact based on the face centered cubic (fcc) packing of balls of 1.7 Å radii. For hard spheres, fcc packing has recently been proved to have the tightest packing.^{34,35} Because the distance between two balls in canonical fcc packing is $2 \times 1.7 = 3.4$ Å, which is greater than the bond length 1.5 Å, we shorten the distance along bonds connecting contacting balls of radius 1.7 Å to 1.5 Å. This mimics the bond length of the model polymer. Unlike fcc packing of hard spheres, bonded monomers here are allowed to have volume overlaps. Additionally, there are some boundary effects because bonds connecting balls in different layer have a distance >1.5 Å. Although mathematically unproven, we conjecture that this artificially constructed polymer represents conformations of SAWs that have very close to maximum compactness.

The packing density of canonical fcc packing by our method is 0.74.¹¹ As described earlier, although fcc packing contains no voids, there are packing crevices or dead spaces that do contribute to the calculation of p_d by our definition.¹¹ In approximately maximum compact polymer, because the distance between bonded balls is shorter than that in fcc packing, p_d can be as high as 0.80 for polymers with a range of chain length.

E. Importance sampling with sequential Monte Carlo

Since we are simulating compact conformations that resemble proteins, we need an efficient method to generate adequate number of conformations satisfying protein-like compactness criteria. Here we use a sequential Monte Carlo (SMC) chain growth strategy,^{36,37} which combines importance sampling and the growth method. The main steps are shown in Fig. 3.

Denote the conformation of a polymer of length *t* as $(x_1, ..., x_t)$ where x_i is the three-dimensional location of the



FIG. 3. The steps of sequential Monte Carlo method applied to improve sampling efficiency.

*i*th monomer. Starting with fixed initial location (x_1, x_2) , we grow polymers by sequentially adding one monomer x_{t+1} to occupy one of the 32-states connecting to the last monomer x_t of the current chain. The monomer x_{t+1} is randomly placed according to а sampling probability $g_{t+1}(x_{t+1}|x_1...x_t)$. In this study, the following function g_{t+1} is used. Let ω be one of the 32-states connected to x_t that satisfies the self-avoiding criterion. First we initialize the number of neighbors $n_{e}(\omega)$ to ω as 1, and the Euclidean distance from ω to the nearest neighbor monomer $d(\omega)$ to 6 Å. We then increment $n_e(\omega)$ by the number of existing monomers within a distance of 6.0 Å to ω . Among these monomers, we identify the monomer x_s that is the nearest neighbor with the shortest Euclidean distance d to ω . We require in addition that the sequence separation |s-t| > 3 so x_s and x_t are not sequence near neighbors. The distance $d(\omega)$ is then replaced by the value of d. The sampling probability is set as

$$g_{t+1}(x_{t+1} = \omega | x_1 \cdots x_t) \propto e^{-E'(\omega)/T'},$$

where $E'(\omega) = \ln\{[d(\omega)]^c/n_e(\omega)\}$ is an artificial "packing energy" favoring more compact conformations, and T' is a pseudotemperature controlling the behavior of sampling. Using this energy function, growth to position ω with close nearest neighbor [small $d(\omega)$] and a large number of neighbors within a 6 Å distance [large $n_e(\omega)$] is favored. Here the adjustable parameter c is used to balance the effect of $d(\omega)$ and $n_e(\omega)$. T' controls the importance of compactness. At low T', conformations generated are compact, but at high T', the compactness criterion becomes less important.

According to the sequential Monte Carlo framework, the importance weight w_{t+1} for the sampled conformation (x_1, \dots, x_{t+1}) is updated as

$$w_{t+1} = w_t \cdot \frac{\pi_{t+1}(x_1 \cdots x_{t+1})}{\pi_t(x_1 \cdots x_t) \cdot g_{t+1}(x_{t+1} | x_1 \cdots x_t)}$$

where $\pi_{t+1}(x_1\cdots x_{t+1})$ is the target distribution at t+1. With a set of weighted samples $\{(x_1^{(j)}, \dots, x_n^{(j)}), w_n^{(j)}\}_{j=1}^m$, statistical inference on the target distribution $\pi_n(x_1, \dots, x_n)$ can be made using

$$E_{\pi_n}[h(x_1,...,x_n)] = \frac{\sum_{j=1}^m w_n^{(j)} \cdot h(x_1^{(j)},...,x_n^{(j)})}{\sum_{j=1}^m w_n^{(j)}}$$
(1)

for most of the proper function h.

F. The target distribution

We wish to generate random samples of polymer with different compactness criterion. This is achieved by using a target distribution π_n which is uniform among all SAWs satisfying a compactness constraint. The constraint is set as follows. First, for each chosen pair values of (T',c), we use the function $e^{-E'(c)/T'}$ to generate 500 random conformations as a trial run. Ignoring the importance weights, we calculate the mean alpha coordination number $z_{\alpha}^*(n,T',c)$ of all the generated conformations. Then we set the target distribution π_n^* as the uniform distribution of all SAWs satisfying $z_{\alpha} \in (0.8 \cdot z_{\alpha}^*(n,T',c), 1.2 \cdot z_{\alpha}^*(n,T',c))$.

We then rerun a large simulation with the same (T',c) parameters and harvest the conformations, using uniform distribution with no restriction on the intermediate target distribution π_t but take the truncated distribution π_n^* as the final target distribution. The truncation is achieved by discarding all generated conformations that does not satisfy the constraint. Typically, the truncation rate is very small (<0.1%). The bias in sampling is fully compensated by proper weighting.

G. Resampling

Because it is easy to have self-avoiding walks to grow into a dead-end, we use resampling to replace dead samples or samples with small weight to improve sampling efficiency.³⁷ Intuitively, we check regularly during the chain growth process whether a particular chain is stuck in a deadend, or is too extended, or has too little weight. If so, this chain is replaced by the replicate of another chain that has the desired compactness. Both duplicate chains will then continue to grow, and the final two surviving chains will be correlated up to the duplication event. Conformations of the monomers added after the duplication will be uncorrelated. This resampling technique targets our simulation to specified configuration space without introducing too much bias, where conformations all have desired compactness (see Fig. 4).

Although we found that the total contact number n_{α} is an excellent parameter for characterizing protein, its calculation involves expensive computation of weighted Delaunay triangulation and alpha shape. We decide to use R_g as a surrogate parameter during resampling. For resampling, we use the empirical relationship $R_g(n) = 2.2 \cdot n^{0.38}$, where *n* is the number of monomers in the polymer, as described in Ref. 38. This relationship has been used as a constraint in NMR pro-



FIG. 4. The steps of the resampling procedure for the sequential Monte Carlo method.

tein structure determination.³⁹ We have the following pseudocode for resampling:³⁷

Procedure RESAMPLING (m, d_s, R_t)

// *m*: Monte Carlo sample size, d_s : steps of looking-back. // R_t : targeting R_{g} .

 $k \leftarrow$ number of dead conformations.

Divide m-k samples randomly into k groups.

for group i=1 to k

Find conformations not picked in previous d_s steps. //Pick the best conformation P_i

 $P_i \leftarrow \text{polymer with } \min |R_g - R_t|$

Replace one of k dead conformations with P_i

Assign both copies of P_i half its original weight.

endfor

Here d_s is used to maintain higher diversity for resampled conformations. That is, conformation that has been picked in the past d_s steps are not available for resampling.

After resampling, the samples with their adjusted weights remain to be *properly weighted* with respect to the original target distribution. We can then calculate the expected alpha coordination number z_{α} , expected packing density p_d , or expected value of any other function *h* using Eq. (1). With these sampling and resampling strategies, we can



FIG. 5. The relationship between packing density p_d , radius of gyration R_g , and the number of residue N in single domain proteins.

successfully grow thousands of self-avoiding walks of chain length up to 2000 using a Linux cluster of 40 CPUs.

III. RESULTS

A. Packing density

Figure 5(a) shows the correlation of packing density p_d with the number of residues *N* in real proteins. Similar relationship has been observed in Ref. 11. Here we further restrict the samples to be of single domain by SCOP annotation.²³ We found that p_d decreases with chain length. That is, short chain proteins have high packing density p_d , but p_d decreases from >0.85 to about 0.74–0.75 when the chain length reaches about 190 residues. After reaching this length, proteins seem to be indifferent about the existence of voids. This suggests that maintaining high packing density is only characteristic of short chain proteins.

B. Radius of gyration of proteins

To identify the factors that dictate the scaling behavior of p_d with residue number N, we need to determine whether such scaling is due to physical constraints of statistical mechanics or the product of extensive optimization by evolution. We study this problem by examining the scaling behavior of p_d with N in random chain polymers generated by computer.

Because of the enormity of conformational space, we focus on random polymers that resemble proteins in some rudimentary sense. One possible criterion is the radius of gyration R_g . This parameter has been widely used as a macroscopic description of protein packing. For single domain proteins, however, we found that there is substantial variance in R_g for proteins of the same chain length [Fig. 5(b)]. Therefore, R_g characterizes protein packing rather poorly, and is unsuitable as a criterion for generating protein-like polymers for our purpose.

C. Alpha contacts

An alternative global description of protein structure is the total number of nonbonded alpha contacts n_{α} defined by the dual simplical complex of the protein. In Fig. 6(a) we plot n_{α} against the total number of atoms in the molecule *n*. As discussed before, these contacts are identified by computing the dual simplicial complex of the molecule.^{11,20} The total number of contacts n_{α} scales linearly with *n*. It also



FIG. 6. The scaling behavior of the total number of nonbonded atomic alpha contacts with the total number of atoms for single domain proteins. Here only contacts from different residues are counted.

scales linearly with the protein chain length (or residue number *N*, data not shown). Regression leads to a linear relationship of $n_{\alpha} = 4.28 \cdot n - 432$, with $R^2 = 0.995$. The alpha contact number n_{α} therefore provides a more accurate global characteristic of protein than radius of gyration R_g . This linear scaling relationship of packing related property is similar to other linear scaling relationships observed for protein, for example, of empirical solvation energy,⁴⁰ protein surface area and protein volume¹¹ with chain length. It is interesting to note that the value of *x*-axis intercept for *n* of the linear regression model suggests that the size of a minimum protein would be in the order of 100 atoms, or about 12–13 residues.

The details of the linear scaling relationship are further examined in Fig. 6(b). It is a replot of Fig. 6(a) after normalization by *n*. It showed that for proteins with 1000 atoms or more (\geq 120 residues), the parameter alpha coordination number $z_{\alpha} = n_{\alpha}/n$ is a constant of about 4.2. For smaller proteins (n < 1,000), z_{α} ranges from 2.5 to 4.0. A nonlinear curve fitting leads to the relationship $z_{\alpha} = a - b/n$, where a= 4.27±0.03, and $b = 4.2 \times 10^2 \pm 26$. We decide to use z_{α} as the criterion to select random polymers generated computationally for packing analysis.

D. Targeted sampling of random chain polymer

Figure 7 shows typical conformations generated with different (T',c) parameters and the conformation of maximally compact polymer. Figure 8 shows the histogram of z_{α} of the conformations at length 2000 without weight adjustment generated using different (T',c) parameters. It can be seen that the histograms for different values of (T',c) do not overlap. This feature demonstrated that with properly chosen (T',c) we can efficiently generate random polymers with z_{α} within a targeted range.



FIG. 7. Examples of self-avoiding walks of length 1000 generated with different sampling probability function $e^{-E'(c)/T'}$ using different (T',c) values. (a) Conformations generated with (T',c)=(1.0, 0.0); (b) (T',c)=(0.67, 0.0); (c) (T',c)=(0.1, 0.6); (d) Approximate maximally compact conformation.



FIG. 8. Self-avoiding walks generated with different (T',c) values do not overlap in z_{α} values. (a) Conformations generated with (T',c) = (1.0,0.0); (b) (T',c) = (0.67,0.0); and (c) (T',c) = (0.1,0.6). Here the numbers of conformations are unweighted. The weighted average z_{α} values are as shown in Fig. 9(a) at length n = 2000.

E. Packing density of random chain polymer

Figure 9(a) shows the relationship of n_{α} associated with each pairs of (T',c) as a function of chain length *n*. It also shows the n_{α} value for the maximally compact conformations. Figure 9(b) shows the relationship of z_{α} and *n*. Note that the targeted z_{α} generated with different (T',c) parameters give rise to different $z_{\alpha} \sim n$ scaling behavior. Because the coarse grained random polymers generated here lack side chains, they are fundamentally different from real proteins. We therefore have experimented with several (T',c) value. We find that protein-like scaling can be obtained for a wide range of (T',c) values.

Figure 9(c) shows the average packing density p_d for all conformations satisfying the constraint specified by different (T',c) values. Except maximally compact conformations, the scaling of $p_d \sim n$ of all other sets of polymers is remarkably similar to that of protein [Fig. 5(a)].

Conformations from set 1 ((T',c) = (1.0,0.0)) are more extended, and have lower average z_{α} [e.g., Fig. 7(a)]. Because there are fewer voids, they also have high p_d . Conformations in set 3 ((T',c) = (0.1,0.6)) are more compact and make more nonbonded contacts and hence have high z_{α} values [e.g., Fig. 7(c)]. They also form more voids, and therefore have lower p_d values. Set 2 ((T',c) = (0.67,0.0)) are conformations whose properties are between those of set 1 and set 3 [e.g., Fig. 7(b)].

The relationship between z_{α} and chain length *n* can be characterized by a nonlinear equation $z_{\alpha} = a - b/n$, similar to that of real proteins. The sets of *a* and *b* obtained by curve fitting are listed in Table I. We emphasize that these ran-

domly generated self-avoiding walks are fundamentally very different from proteins: all residues are of uniform size, there are no side chains, and there is no hydrophobic or any other type of physical interactions in these polymers. Because it is impossible to quantitatively define a similarity metric that measures how different these polymers are from proteins, we are not able to decide which specific values of (T',c) are optimal for modeling protein packing. Nevertheless, the scaling of $p_d \sim n$ for all (T',c) values is qualitatively quite similar to that of real proteins.

The relationship between p_d and z_α at chain length 1800 for self-avoiding walks generated with different parameters (T',c) are shown in Fig. 10. The p_d of both extended and maximally compact conformations have high p_d values, but conformations with an intermediate value of z_α contain voids and have smaller p_d values. This is similar to the relationship of p_d and a compactness parameter ρ (equivalent to z_α used here) studied in two-dimensional lattice (see Fig. 8 in Ref. 37).

IV. SUMMARY AND DISCUSSION

It is well acknowledged that protein has high packing density p_d , as high as that of crystalline solids.⁸ However, recent study suggested that there are numerous voids and pockets in proteins.¹¹ It was also found that about 1/3 of the residues in a protein deviates from the fcc close packing and have random positions.⁵ The simulation results presented here indicate that chain connectivity, excluded volume, and global compactness are the main determinants of the scaling behavior of voids and chain length in proteins. Unlike maxi-



FIG. 9. The relationship of alpha contact number n_{α} , alpha coordination number z_{α} , and the packing density p_d of random compact and maximally compact self-avoiding walks. The curves are sampled following the function $e^{-E'(c)/T'}$, with different T' and c values. Each data point is an average of 10 runs of sample size 600.

mally compact polymers which maintains high packing density at all chain length, proteins and simple near compact polymers have large p_d values only for relatively short chains. When the chain length reaches 190 residues for protein and about 600–700 for chain polymers, proteins and polymers have lower packing density and are quite tolerant to the formation of voids.

The global compactness is a necessary condition for the observed protein-like scaling behavior. However, not all parameters related to voids and compactness are equally appropriate. The data shown in Fig. 6 suggests that the alpha coordination number z_{α} reflects basic intrinsic compactness properties of protein, which is absent in the widely used parameter R_g , the radius of gyration. The advantage of parameters such as z_{α} emphasizes the importance of accurate description of protein geometry and structure.

The parameters p_d is biased towards short chain proteins. By definition, a polymer formed by 2, 3 or a small number of monomers do not have long enough chains to form voids, therefore all will have $p_d=1.0$. When chain length becomes longer, voids appear. Similarly, z_{α} is also

TABLE I. The relationship between z_{α} and *n* can be described by the equation $z_{\alpha} = a - b/n$. The estimated values of *a* and *b*, along with standard deviations in parentheses, are listed for three different sets of conformations generated with different (T', c) parameters which characterize the sampling probability. The values of *a* and *b* for real proteins are also listed.

T'	С	а	b
1.0	0.0	1.88(0.02)	72(7)
0.67	0.0	2.24(0.02)	67(6)
0.1	0.6	3.68(0.02)	93(7)
Native proteins		4.27(0.04)	423(27)

biased towards short chains. For very short chain polymers where the chain has few turns, few nonbonded contacts exist and no voids are formed. In this case, z_{α} is low and p_d is high. However, this small size effect disappears rapidly for our model conformations for maximally compact polymers. Small size effect therefore does not fully account for the scaling behavior of p_d and z_{α} in proteins and in simulated random polymers.

There are major differences between self-avoiding walks we generated and real protein structures. Our SAWs have no side chains, and belong to the coarse-grain model where one monomer is represented as a ball. In addition, the target distribution of SMC sampling is the truncated uniform distribution of all geometrically feasible conformations. The trunca-



FIG. 10. The relationship of p_d and $z_{\alpha} = n_c/n$ for self-avoiding walks generated using different (T', c) parameters at chain length 1800.

In this paper, we describe a novel approach to overcome the attrition problem in generating long chain compact selfavoiding walk. With sequential importance sampling and resampling, we have developed an algorithm that effectively sample rare events, i.e., compact self-avoiding walks. Success in generating thousands of off-lattice self-avoiding walks satisfying various desired compactness requirement is essential for studying the scaling behavior of p_d in random off-lattice self-avoiding walks.

The main result of this paper is that with sequential Monte Carlo techniques it is not difficult to reproduce protein-like scaling behavior of packing density p_d and chain length n in generic chain polymers. With the guidance of rudimentary requirement of z_{α} , this can be achieved under a wide range of $z_{\alpha} \sim n$ relationships. We therefore conclude that proteins retain the same packing property of generic compact chain polymers. We further conclude that proteins are unlikely to be optimized by evolution to eliminated packing voids. This is in support of the insightful comments of Richards who suggested that an appropriate level of underpacking would be important for evolution to occur through random mutations.¹⁰

Our study showed the importance of generic geometric packing related to z_{α} in reproducing protein-like $p_d \sim n$ scaling behavior. To test further the role of geometric packing, the next step would be to examine the $p_d \sim n$ scaling by generating more realistic compact random polymers with perhaps monomers of different sizes to model the side chain effects. Furthermore, with sequential Monte Carlo and other advanced sampling methods, various models of explicit side chains can be attached to main chain monomers. In addition, one could introduce various alphabet sets for the residues (such as the HP model) and corresponding potential energy function H. In this case, the target distribution can be the Boltzmann distribution $\pi \propto \exp(-H/T)$ instead of the uniform distribution of all SAWs. It would also be interesting to examine the $p_d \sim n$ scaling of polymers of random sequences and of protein-like sequences with low energy in compact states.

ACKNOWLEDGMENTS

The authors thank Professor Herbert Edelsbrunner and Professor Luhua Lai for valuable discussions. This work is supported by funding from the National Science Foundation CAREER DBI0133856, DBI0078270, DMS0073601, CCR9980599 and American Chemical Society/Petroleum Research Fund 35616-G7.

- ¹C. Chothia, M. Levitt, and D. Richardson, J. Mol. Biol. 145, 214 (1981).
- ²C. Chothia and A. V. Finkelstein, Annu. Rev. Biochem. **59**, 1007 (1990).
 ³A. Maritan, C. Micheletti, A. Trovato, and J. R. Banavar, Nature (London) **406**, 287 (2000).
- ⁴J. R. Banavar, A. Maritan, C. Micheletti, and A. Trovato, Proteins 47, 315 (2002).
- ⁵Z. Bagci, R. L. Jernigan, and I. Bahar, J. Chem. Phys. **116**, 2269 (2002).
- ⁶P. B. Sigler, Z. Xu, H. S. Rye, S. G. Burston, W. A. Fenton, and A. L.
- Horwich, Annu. Rev. Biochem. 67, 581 (1998).
- ⁷D. A. Doyle, J. Morais Cabral, R. A. Pfuetzner, A. Kuo, J. M. Gulbis, S. L. Cohen, B. T. Chait, and R. MacKinnon, Science **280**, 69 (1998).
- ⁸F. M. Richards and W. A. Lim, Q. Rev. Biophys. 26, 423 (1994).
- ⁹J. Hermans and H. A. Scheraga, J. Am. Chem. Soc. 83, 3283 (1961).
- ¹⁰F. M. Richards, Cell. Mol. Life Sci. 53, 790 (1997).
- ¹¹ J. Liang and K. A. Dill, Biophys. J. 81, 751 (2001).
- ¹² A. E. Eriksson, W. A. Baase, X. J. Zhang, D. W. Heinz, M. Blaber, E. P. Baldwin, and B. W. Matthews, Science **255**, 178 (1992).
- ¹³P. L. Privalov, J. Mol. Biol. **258**, 707 (1996).
- ¹⁴B. I. Dahiyat and S. L. Mayo, Science 221, 709 (1997).
- ¹⁵O. B. Ptitsyn, J. Mol. Biol. **278**, 655 (1998).
- ¹⁶O. B. Ptitsyn and K. L. Ting, J. Mol. Biol. **291**, 671 (1999).
- ¹⁷L. Mirny and E. Shakhnovich, Annu. Rev. Biophys. Biomol. Struct. **30**, 361 (2001).
- ¹⁸ H. Li, R. Helling, C. Tang, and N. Wingreen, Science **273**, 666 (1996).
 ¹⁹ R. Mélin, H. Li, N. Wingreen, and C. Tang, J. Chem. Phys. **110**, 1252 (1999).
- ²⁰ J. Liang, H. Edelsbrunner, P. Fu, P. V. Sudhakar, and S. Subramaniam, Proteins 33, 1 (1998).
- ²¹ J. Liang, H. Edelsbrunner, P. Fu, P. V. Sudhakar, and S. Subramaniam, Proteins **33**, 18 (1998).
- ²²U. Hobohm and C. Sander, Protein Sci. 3, 522 (1994).
- ²³ A. G. Murzin, S. E. Brenner, T. Hubbard, and C. Chothia, J. Mol. Biol. 247, 536 (1995).
- ²⁴ H. Edelsbrunner, M. Facello, P. Fu, and J. Liang, in *Proceedings of the 28th Annual Hawaii International Conference System Sciences*, Los Alamitos, California (IEEE Computer Society, California, 1995).
- ²⁵K. P. Peters, J. Fauck, and C. Frömmel, J. Mol. Biol. **256**, 201 (1996).
- ²⁶J. Liang, H. Edelsbrunner, and C. Woodward, Protein Sci. 7, 1884 (1998).
- ²⁷ J. Liang and S. Subranmaniam, Biophys. J. **73**, 1830 (1997).
- ²⁸ J. Liang and M. P. McGee, Biophys. J. **75**, 573 (1998).
- ²⁹H. Edelsbrunner, Discrete Comput. Geom. 13, 415 (1995).
- ³⁰H. Edeslbrunner, M. Facello, and J. Liang, Disc. Appl. Math. **88**, 18 (1998).
- ³¹H. Edelsbrunner and E. Mücke, ACM Trans. Graphics 13, 43 (1994).
- ³²B. H. Park and M. Levitt, J. Mol. Biol. **249**, 493 (1995).
- ³³M. N. Rosenbluth and A. W. Rosenbluth, J. Chem. Phys. 23, 356 (1955).
- ³⁴T. C. Hale, The Kepler conjecture, http://www.math.pitt.edu/thales
- ³⁵N. J. A. Sloane, Nature (London) **395**, 435 (1998).
- ³⁶J. S. Liu and R. Chen, J. Am. Stat. Assoc. **93**, 1032 (1998).
- ³⁷J. Liang, J. Zhang, and R. Chen, J. Chem. Phys. **117**, 3511 (2002).
- ³⁸J. Skolnick, A. Kolinski, and A. R. Ortiz, J. Mol. Biol. 265, 217 (1997).
- ³⁹X. Huang and R. Powers, J. Am. Chem. Soc. **123**, 3834 (2001).
- ⁴⁰L. Chiche, L. M. Gregoret, F. E. Cohen, and P. A. Kollman, Proc. Natl. Acad. Sci. U.S.A. **87**, 3240 (1990).